| Chermoi | oyn amic Functio | ns of Formation | FOR $Pb(CH_3)$, | | | | | | | |
|--------------------|-------------------------------------|------------------------------|----------------------------------|--|--|--|--|--|--|--|
| | (g) Ani | $D Sn(CH_3)_4(g)$ | | | | | | | | |
| Τ , °K. | $(\Delta H^0)_{\rm f}$, kcal./mole | (ΔF^0) f, kcal./mole | $(\Delta S^0)_{\rm f}/{ m mole}$ | | | | | | | |
| $Pb(CH_3)_4(g)$ | | | | | | | | | | |
| 300 | 3.2 ± 3 | 35.6 ± 3.6 | -108 ± 2 | | | | | | | |
| 400 | 1.2 | 46.0 | -112 | | | | | | | |
| 500 | -0.5 | 57.0 | -115 | | | | | | | |
| 600 | -1.9 | 68.1 | -117 | | | | | | | |
| | Sn(| $(CH_3)_4(g)^a$ | | | | | | | | |
| 3 00 | -13.6 ± 10 | 19.0 ± 10 | -109 | | | | | | | |
| 400 | -15.7 | 30.2 | -115 | | | | | | | |
| 500 | -17.5 | 41.9 | -119 | | | | | | | |
| ^a Based | 1 on $(\Delta H^0)_{\rm f}$ for St | $nO_2 = -138.8 \text{ kc}$ | al./mole. | | | | | | | |

energy and entropy of formation of SnO₂ given by N.B.S.,¹⁸ the heat of formation, (ΔH^0) , at 300°K.

(18) N.B.S. Circular 500 "Selected Values of Chemical Thermodynamic Properties." was calculated. Using this, the calculated thermodynamic functions for Sn(s) and Sn(CH₃)₄(g), and the functions for C(graphite) and H₂(g) given in N.B.S. tables, a value of E_0^0 of -5.5 ± 10 kcal./ mole, was obtained. The thermodynamic functions for Sn(s) were obtained from data given by Kelley¹⁹ and Klinkhardt¹³ and are tabulated in Table XII. From E_0^0 the functions $(\Delta H^0)_{\rm f}$, $(\Delta F^0)_{\rm f}$ and $(\Delta S^0)_{\rm f}$ were calculated by standard methods and are tabulated in Table XIII from 300 to 500°K.

The authors wish to acknowledge financial support from the Office of Naval Research Contract N 8 onr-72700. The authors wish to express their gratitude to Professor Richard C. Lord, of the Massachusetts Institute of Technology, for the use of the equipment in the Spectroscopy Laboratory.

(19) K. K. Kelley, Bureau of Mines Bulletin 350, (1932).

Manhattan, Kansas Storrs, Conn.

[FROM THE UNIVERSITY OF CALIFORNIA, DEPARTMENT OF CHEMISTRY AND CHEMICAL ENGINEERING, AND THE RADIATION LABORATORY, BERKELEY]

Absolute Entropies in Liquid Ammonia

By Wendell M. Latimer and William L. Jolly

RECEIVED MAY 15, 1953

It is shown that a simple relationship exists between the partial molal entropies of ions in water and liquid ammonia if the absolute entropy of H^+ is taken as -2 e.u. in water and -25 e.u. in ammonia. A new tabulation of thermodynamic data for ions in liquid ammonia is given.

(1951).

From studies of relative ionic entropies in methanol,¹ and ammonia^{2,3} and from the entropies of

ionization of methanol, ethanol, acetic acid and ammonia,4 it has been concluded that the absolute ionic entropies in these solvents are considerably more negative than the corresponding absolute entropies in water. The arguments for this conclusion are: (1) The entropies of solution of electrolytes in water are generally more positive than in the non-aqueous solvents, and (2) the entropy of ionization of water is more positive than the corresponding entropies of ionization of the non-aqueous solvents.

In the case of water, the absolute entropy of the hydrogen ion has been determined as -2.1 e.u.^5 (very close to the usual arbitrary value of 0 e.u.)

(1) W. M. Latimer and C. M. Slansky,

THIS JOURNAL, 62, 2019 (1940).

(2) W. L. Jolly, Chem. Revs., 50, 351 (1952).

(3) S. P. Wolsky, E. J. Zdanuk and L. V. Coulter, THIS JOURNAL, 74, 6196 (1952).

(4) W. L. Jolly, *ibid.*, **74**, 6199 (1952).

(5) J. C. Goodrich, F. M. Goyan, E. E. Morse, R. G. Preston and M. B. Young, *ibid.*, **72**, 4411 (1950).

and monatomic ion entropies may be well represented by the equation of Powell and Latimer⁶





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|------------------------------------|--------------------------------------------|--------------------------------------------|-------------|------------------|--------------------------------------------|--------------------------------------------|-------------|--|--|
| THERMODYNAMIC FUNCTIONS AT 298 °K. | | | | | | | | | |
| | $\Delta H_{\rm f}^{\circ}$, kcal./mole | $\Delta F_{\rm f}^{\circ}$, kcal./mole | S°, e.u. | | $\Delta H_{\rm f}^{\circ}$, kcal./mole | $\Delta F_{\rm f}^{\circ}$, kcal./mole | S°, e.u. | | |
| H+ | 0 | 0 | 0 | SCN- | -11.8 | | | | |
| H₂O | -71.6 | | | Pb++ | 21.0 | 13 | 11 | | |
| F- | | -80.8 | | Pb9 | 64 | | | | |
| C1- | -65.7 | -44.1 | -30.3 | T1+ | | 5,7 | | | |
| C1O3 ⁻ | -47.7 | -17.7 | 15 | Zn++ | | -25 | | | |
| Br - | -59.0 | -39.9 | -30.3 | Hg++ | 45.2 | 31 | 35 | | |
| I - | -45.3 | -29 | -25 | Ag+ | 26 | 17.5 | 23 | | |
| IO3- | | -42.4 | | Mn++, | | -26 | | | |
| Se ₂ - | -19.0 | | | BH_4^- | -15.7 | | | | |
| Te ₂ | -27.8 | | | Ca++ | -100.0 | -100 | -21 | | |
| Te ₄ | -29.7 | | | Li ⁺⁺ | -49 | -54 | 8 | | |
| NO3- | -77.6 | -42.8 | - 5 | Na + | -38.1 | -43.6 | 15.1 | | |
| NH3 | -16.1 | -2.7 | 24.7 | K + | -40.5 | -47.0 | 21.4 | | |
| NH₄+ | -16.1 | -2.7 | 24.7 | Rb+ | -39 | -47.5 | 29 | | |
| NH2 ⁻ | 10.1 | 34 | -10 | Cs+ | -40.5 | -48 | 29 | | |

 $1/2e_2$ (am)

e⁻(am)

both the positive and negative ions fall upon a line if the absolute entropy of H^+ in ammonia is taken as -25 e.u. This plot is shown in Fig. 1. Since the slope of the line (as drawn) is unity, the coefficient of the Z/r_e^2 term in the Powell-Latimer equation must be approximately the same for both ammonia and water. At least for monatomic ions, it would appear that absolute entropies in water are simply 23 units higher than in ammonia.

-24.6

-20.0

Values for NO_3^- and ClO_3^- are included in Fig. 1, but since the Powell-Latimer equation applies only to the monatomic ions, it is not surprising that they fall off the curve. The experimental data for dipositive ions in ammonia are so incomplete that it is not possible to ascertain the dependence of the entropies on the charge on the ion with any accuracy and the tentative values for Ca^{++} , Pb^{++} and Hg^{++} have not been plotted in Fig. 1.

40.5

43.5

44.4

A recalculation has been made of many of the thermodynamic functions for ions in ammonia.⁷ A number of these new values differ from those previously tabulated by Jolly,² and a new summary is therefore given in Table I.

(7) For the detailed calculations see W. L. Jolly, University of California Radiation Laboratory Report No. UCRL-2201, May, 1953. BERKELEY, CALIF.

[CONTRIBUTION FROM THE NOVES CHEMICAL LABORATORY, UNIVERSITY OF ILLINOIS]

Verification of a Theory of Irreversible Polarographic Waves

BY PEKKA KIVALO, K. B. OLDHAM AND H. A. LAITINEN

RECEIVED APRIL 30, 1953

Two independent schools have recently derived equations of totally irreversible waves; the equations are identical. The equation predicts that at the foot of the wave, the current obeys relationships that are typical of kinetic control; at the head of the wave, however, diffusion controlled relationships apply. Several expressions have been derived from the basic equation. Experimental verification of these expressions furnishes a convincing proof of the validity of the theory of irreversible waves.

During recent years, two independent approaches have been made to the problem of the interpretation of completely irreversible polarographic waves. The treatment employed by Delahay and Strassner¹⁻³ differs only in minor respects from that due to Evans, Hush and Oldham.⁴⁻⁶ The basic equation of the polarographic wave, as derived by Delahay,^{1,2} is identical in all but mathematical terminology from that obtained by Hush and Oldham.⁶ The symbolism and method of the latter workers has been retained in the present paper because it

- (3) P. Delahay, ONR report No. 7, Project NR-051-258 (1952).
 (4) M. G. Evans and N. S. Hush, J. chim. phys., 49, C159 (1952).
- (5) N. S. Hush and K. B. Oldham, Nature, in press.
- (6) N. S. Hush and K. B. Oldham, Trans. Faraday Soc., in press.

is more amenable to mathematical treatment than the graphical procedure of Delahay.

The purpose of the theoretical section of the present paper is to present some of the implications of the theory and, in particular, to derive relationships that may be readily tested experimentally. The experimental section is devoted to the verification of these relationships. Such verifications of the theory as have hitherto been reported, ^{2.6} have been of a somewhat scanty nature.

Implications of the Equation of the Irreversible Wave.—The basic equation derived by Hush and Oldham⁶ to express the instantaneous current, i_t (in microamperes), passed by a dropping mercury electrode in consequence of the totally irreversible reduction

 $\mathbf{2}$

NH₂OH

Guanidinium ion

⁽¹⁾ P. Delahay, THIS JOURNAL, 73, 4944 (1951).

⁽²⁾ P. Delahay and J. E. Strassner, ibid., 73, 5219 (1951).